

## Integrals formulas involving confluent hypergeometric Functions of three variables $\Phi_2^{(3)}$ and $\Psi_2^{(3)}$

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### Abstract

The aim of this paper is to establish two general integral formulas involving confluent hypergeometric functions of three variables  $\Phi_2^{(3)}$  and  $\Psi_2^{(3)}$  with the help of two extension formulas for Lauricella's functions of three variables  $F_A^{(3)}$  and  $F_D^{(3)}$  due to Atash [1] and Atash and Bellehaj [2]. Some applications of our main results are also presented.

**Keywords:** Integral formulas, Hypergeometric functions, Dixon's theorem, Kummer's theorem

### 1. Introduction

The Lauricella's functions of three variables  $F_A^{(3)}$  and  $F_D^{(3)}$  are defined and represented as follows [7]:

$$F_A^{(3)}(a, b_1, b_2, b_3; c_1, c_2, c_3; x_1, x_2, x_3) = \sum_{m_1, m_2, m_3=0}^{\infty} \frac{(a)_{m_1+m_2+m_3} (b_1)_{m_1} (b_2)_{m_2} (b_3)_{m_3}}{(c_1)_{m_1} (c_2)_{m_2} (c_3)_{m_3}} \frac{x_1^{m_1} x_2^{m_2} x_3^{m_3}}{m_1! m_2! m_3!} \quad (1.1)$$

$$|x_1| + |x_2| + |x_3| < 1$$

and

$$F_D^{(3)}(a, b_1, b_2, b_3; d; x_1, x_2, x_3) = \sum_{m_1, m_2, m_3=0}^{\infty} \frac{(a)_{m_1+m_2+m_3} (b_1)_{m_1} (b_2)_{m_2} (b_3)_{m_3}}{(d)_{m_1+m_2+m_3}} \frac{x_1^{m_1} x_2^{m_2} x_3^{m_3}}{m_1! m_2! m_3!} \quad (1.2)$$

$$\max\{|x_1|, |x_2|, |x_3|\} < 1,$$

where  $(a)_n$  is the Pochhammer's symbol defined by

$$(a)_n = \begin{cases} 1 & , \text{ if } n = 0 \\ a(a+1)(a+2)\dots(a+n-1) & , \text{ if } n = 1,2,3,\dots \end{cases} \quad (1.3)$$

The Laplace integral representations of Lauricella functions  $F_A^{(3)}$  and  $F_D^{(3)}$  are given by Exton (see [3]):

$$F_A^{(3)}(a, b_1, b_2, b_3; c_1, c_2, c_3; x_1, x_2, x_3) = \frac{1}{\Gamma(b_1)\Gamma(b_2)\Gamma(b_3)} \int_0^{\infty} \int_0^{\infty} \int_0^{\infty} e^{-t_1-t_2-t_3} t_1^{b_1-1} t_2^{b_2-1} t_3^{b_3-1} \times \Psi_2^{(3)}(a; c_1, c_2, c_3; x_1 t_1, x_2 t_2, x_3 t_3) dt_1 dt_2 dt_3, \quad (1.4)$$

$$\text{Re}(b_1), \text{Re}(b_2), \text{Re}(b_3) > 0$$

and

$$F_D^{(3)}(a, b_1, b_2, b_3; d; x_1, x_2, x_3) = \frac{1}{\Gamma(a)} \int_0^\infty e^{-s} s^{a-1} \Phi_2^{(3)}(b_1, b_2, b_3; d; x_1 s, x_2 s, x_3 s) ds, \tag{1.5}$$

$\text{Re}(a) > 0,$

where the functions  $\Phi_2^{(3)}$  and  $\Psi_2^{(3)}$  are the confluent hypergeometric functions defined by Srivastava[7] .

$$\Phi_2^{(3)}(b_1, b_2, b_3; c; x_1, x_2, x_3) = \sum_{m_1, m_2, m_3=0}^\infty \frac{(b_1)_{m_1} (b_2)_{m_2} (b_3)_{m_3}}{(c)_{m_1+m_2+m_3}} \frac{x_1^{m_1} x_2^{m_2} x_3^{m_3}}{m_1! m_2! m_3!} \tag{1.6}$$

and

$$\Psi_2^{(3)}(a; b_1, b_2, b_3; x_1, x_2, x_3) = \sum_{m_1, m_2, m_3=0}^\infty \frac{(a)_{m_1+m_2+m_3}}{(b_1)_{m_1} (b_2)_{m_2} (b_3)_{m_3}} \frac{x_1^{m_1} x_2^{m_2} x_3^{m_3}}{m_1! m_2! m_3!} . \tag{1.7}$$

The Exton's double hypergeometric functions are defined by [4]

$$X \begin{matrix} A: B; B' \\ C: D; D' \end{matrix} \left[ \begin{matrix} (a):(b); (b') \\ (c):(d); (d') \end{matrix} ; x, y \right] = \sum_{m,n=0}^\infty \frac{((a))_{2m+n} ((b))_m ((b'))_n x^m y^n}{((c))_{2m+n} ((d))_m ((d'))_n m! n!}, \tag{1.8}$$

where the symbol  $((a))_m$  denotes the product  $\prod_{j=1}^A (a_j)_m$  .

## 2.MainIntegralFormulas

By employing the generalized Dixon's theorem [5] and the generalized Kummer's theorem [6],Atash [1] and Atash and Bellehaj [2] derived the following two extension formulas for Lauricella's functions of three variables  $F_A^{(3)}$  and  $F_D^{(3)}$  :

$$\begin{aligned} & F_A^{(3)}(a, a', b-i, b; d, c, c+i+j; x, y, -y) \\ &= \sum_{m=0}^\infty \sum_{n=0}^\infty \frac{(a)_{m+2n} (a')_m (b-i)_{2n} x^m (-y^2/4)^n}{(d)_m (c)_{2n} m! n!} \\ & \times \frac{2^{2(2n+c-1)+i+j} \Gamma(1-b+i-2n) \Gamma(c+i+j) \Gamma(b-\frac{1}{2}|i|-\frac{1}{2}i) \Gamma(1-2n-c-\frac{1}{2}(i+j+|i+j|)}{\Gamma(b)\Gamma(1-c-2n)\Gamma(2c-1+i+j+2n)\Gamma(c-b+i+j)} \\ & \times C_{i,j} \frac{\Gamma(n+c-\frac{1}{2}+[\frac{i+j+1}{2}])\Gamma(n-b+c+i+[\frac{j+1}{2}])}{\Gamma(\frac{1}{2})\Gamma(1-b-n+[\frac{i}{2}])} \\ & - \frac{ay}{2} \sum_{m=0}^\infty \sum_{n=0}^\infty \frac{(a+1)_{m+2n} (a')_m (b-i)_{2n+1} x^m (-y^2/4)^n}{(d)_m (c)_{2n+1} m! n!} \\ & \times \frac{2^{2(c+2n)+i+j} \Gamma(i-b-2n) \Gamma(c+i+j) \Gamma(b-\frac{1}{2}|i|-\frac{1}{2}i) \Gamma(-2n-c-\frac{1}{2}(i+j+|i+j|)}{\Gamma(b)\Gamma(-2n-c)\Gamma(2c+i+j+2n)\Gamma(c-b+i+j)} \\ & \times D_{i,j} \frac{\Gamma(\frac{1}{2}+n+c+[\frac{i+j}{2}])\Gamma(1+n-b+c+i+[\frac{j}{2}])}{\Gamma(\frac{1}{2})\Gamma(-n-b+[\frac{i+1}{2}])}, \tag{2.1} \end{aligned}$$

for  $i = -3, -2, -1, 0, 1, 2$  and  $j = 0, 1, 2, 3$

and

$$F_D^{(3)}(a, b-i, b, c; d; x, -x, y)$$

$$\begin{aligned}
 &= \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} \frac{(a)_{2m+n} (b-i)_{2m} (c)_n x^{2m} y^n}{(d)_{2m+n} (2m)! n!} \\
 &\left\{ A'_i \frac{2^{2m} \Gamma(\frac{1}{2}) \Gamma(1-2m-b+i) \Gamma(1-b)}{\Gamma(1-b+\frac{1}{2}(i+|i|)) \Gamma(-m+\frac{1}{2}i+\frac{1}{2}-[\frac{1+i}{2}]) \Gamma(1-m-b+\frac{1}{2}i)} \right. \\
 &+ B'_i \frac{2^{2m} \Gamma(\frac{1}{2}) \Gamma(1-2m-b+i) \Gamma(1-b)}{\Gamma(1-b+\frac{1}{2}(i+|i|)) \Gamma(-m+\frac{1}{2}i-[\frac{i}{2}]) \Gamma(-m+\frac{1}{2}-b+\frac{1}{2}i)} \left. \right\} \\
 &+ \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} \frac{(a)_{2m+n+1} (b-i)_{2m+1} (c)_n x^{2m+1} y^n}{(d)_{2m+n+1} (2m+1)! n!} \\
 &\times \left\{ A''_i \frac{2^{2m+1} \Gamma(\frac{1}{2}) \Gamma(-2m-b+i) \Gamma(1-b)}{\Gamma(1-b+\frac{1}{2}(i+|i|)) \Gamma(-m+\frac{1}{2}i-[\frac{1+i}{2}]) \Gamma(-m+\frac{1}{2}-b+\frac{1}{2}i)} \right. \\
 &+ B''_i \frac{2^{2m+1} \Gamma(\frac{1}{2}) \Gamma(-2m-b+i) \Gamma(1-b)}{\Gamma(1-b+\frac{1}{2}(i+|i|)) \Gamma(-m-\frac{1}{2}+\frac{1}{2}i-[\frac{i}{2}]) \Gamma(-m-b+\frac{1}{2}i)} \left. \right\} \quad (2.2) \\
 &\text{for } i=0, \pm 1, \pm 2, \pm 3, \pm 4, \pm 5,
 \end{aligned}$$

where  $[x]$  denotes the greatest integer less than or equal to  $x$  and  $|x|$  denotes the usual absolute value of  $x$ . The coefficient  $C_{i,j}$  can be obtained from the table of  $A_{i,j}$  given in [5] by replacing  $a$  and  $c$  by  $-2n$  and  $1-c-2n$  and the coefficient  $D_{i,j}$  can be obtained from the table of  $B_{i,j}$  given in [5] by replacing  $a$  and  $c$  by  $-2n-1$  and  $-c-2n$  respectively.

The coefficients  $A'_i$  and  $B'_i$  can be obtained from the tables of  $A_i$  and  $B_i$  given in [6] by taking  $a = -2m$  and the coefficients  $A''_i$  and  $B''_i$  can be obtained from the same tables of  $A_i$  and  $B_i$  by taking  $a = -2m-1$ .

In (1.4) replacing  $b_1, b_2, b_3, c_1, c_2, c_3, x_1, x_2$  and  $x_3$  by  $a', b-i, b, d, c, c+i+j, x, y$  and  $-y$  respectively and using the result (2.1), we get the following general integral involving confluent hypergeometric function of three variables  $\Psi_2^{(3)}$ :

**First Integral**

$$\begin{aligned}
 &\frac{1}{\Gamma(a')\Gamma(b-i)\Gamma(b)} \int_0^{\infty} \int_0^{\infty} \int_0^{\infty} e^{-t_1-t_2-t_3} t_1^{a'-1} t_2^{b-i-1} t_3^{b-1} \\
 &\times \Psi_2^{(3)}(a; d, c, c+i+j; xt_1, yt_2, -yt_3) dt_1 dt_2 dt_3 . \\
 &= \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} \frac{(a)_{m+2n} (a')_m (b-i)_{2n} x^m (-y^2/4)^n}{(d)_m (c)_{2n} m! n!} \\
 &\times \frac{2^{2(2n+c-1)+i+j} \Gamma(1-b+i-2n) \Gamma(c+i+j) \Gamma(b-\frac{1}{2}|i|-\frac{1}{2}i) \Gamma(1-2n-c-\frac{1}{2}(i+j+|i+j|)}{\Gamma(b)\Gamma(1-c-2n)\Gamma(2c-1+i+j+2n)\Gamma(c-b+i+j)} \\
 &\times C_{i,j} \frac{\Gamma(n+c-\frac{1}{2}+[\frac{i+j+1}{2}]) \Gamma(n-b+c+i+[\frac{j+1}{2}])}{\Gamma(\frac{1}{2}) \Gamma(1-b-n+[\frac{i}{2}])} \\
 &- \frac{ay}{2} \sum_{m=0}^{\infty} \sum_{n=0}^{\infty} \frac{(a+1)_{m+2n} (a')_m (b-i)_{2n+1} x^m (-y^2/4)^n}{(d)_m (c)_{2n+1} m! n!}
 \end{aligned}$$

$$\begin{aligned} & \times \frac{2^{2(c+2n)+i+j} \Gamma(i-b-2n) \Gamma(c+i+j) \Gamma(b-\frac{1}{2}|i|-\frac{1}{2}i) \Gamma(-2n-c-\frac{1}{2}(i+j+|i+j|))}{\Gamma(b) \Gamma(-2n-c) \Gamma(2c+i+j+2n) \Gamma(c-b+i+j)} \\ & \times D_{i,j} \frac{\Gamma(\frac{1}{2}+n+c+\lceil \frac{i+j}{2} \rceil) \Gamma(1+n-b+c+i+\lceil \frac{j}{2} \rceil)}{\Gamma(\frac{1}{2}) \Gamma(-n-b+\lceil \frac{i+1}{2} \rceil)}, \quad (2.3) \\ & \text{for } i = -3, -2, -1, 0, 1, 2 \text{ and } j = 0, 1, 2, 3. \end{aligned}$$

Next, in (1.5), replacing  $b_1, b_2, b_3, x_1, x_2$  and  $x_3$  by  $b-i, b, c, x, -x$  and  $y$  respectively and using the result (2.2), we get the following general integral involving confluent hypergeometric function of three variables  $\Phi_2^{(3)}$  :

**Second Integral**

$$\begin{aligned} & \frac{1}{\Gamma(a)} \int_0^\infty e^{-s} s^{a-1} \Phi_2^{(3)}(b-i, b, c; d; xs, -xs, ys) ds \\ & = \sum_{m=0}^\infty \sum_{n=0}^\infty \frac{(a)_{2m+n} (b-i)_{2m} (c)_n x^{2m} y^n}{(d)_{2m+n} (2m)! n!} \\ & \left\{ A'_i \frac{2^{2m} \Gamma(\frac{1}{2}) \Gamma(1-2m-b+i) \Gamma(1-b)}{\Gamma(1-b+\frac{1}{2}(i+|i|)) \Gamma(-m+\frac{1}{2}i+\frac{1}{2}-\lceil \frac{1+i}{2} \rceil) \Gamma(1-m-b+\frac{1}{2}i)} \right. \\ & \left. + B'_i \frac{2^{2m} \Gamma(\frac{1}{2}) \Gamma(1-2m-b+i) \Gamma(1-b)}{\Gamma(1-b+\frac{1}{2}(i+|i|)) \Gamma(-m+\frac{1}{2}i-\lceil \frac{i}{2} \rceil) \Gamma(-m+\frac{1}{2}-b+\frac{1}{2}i)} \right\} \\ & + \sum_{m=0}^\infty \sum_{n=0}^\infty \frac{(a)_{2m+n+1} (b-i)_{2m+1} (c)_n x^{2m+1} y^n}{(d)_{2m+n+1} (2m+1)! n!} \\ & \times \left\{ A''_i \frac{2^{2m+1} \Gamma(\frac{1}{2}) \Gamma(-2m-b+i) \Gamma(1-b)}{\Gamma(1-b+\frac{1}{2}(i+|i|)) \Gamma(-m+\frac{1}{2}i-\lceil \frac{1+i}{2} \rceil) \Gamma(-m+\frac{1}{2}-b+\frac{1}{2}i)} \right. \\ & \left. + B''_i \frac{2^{2m+1} \Gamma(\frac{1}{2}) \Gamma(-2m-b+i) \Gamma(1-b)}{\Gamma(1-b+\frac{1}{2}(i+|i|)) \Gamma(-m-\frac{1}{2}+\frac{1}{2}i-\lceil \frac{i}{2} \rceil) \Gamma(-m-b+\frac{1}{2}i)} \right\} \quad (2.4) \\ & \text{for } i = 0, \pm 1, \pm 2, \pm 3, \pm 4, \pm 5. \end{aligned}$$

**3. Applications**

In this section, we will use in each case the following results [7]:

$$(a)_n = \frac{\Gamma(a+n)}{\Gamma(a)}, \quad a \neq 0, -1, -2, \dots \quad (3.1)$$

$$(a)_{2n} = 2^{2n} \left(\frac{1}{2}a\right)_n \left(\frac{1}{2}a + \frac{1}{2}\right)_n \quad (3.2)$$

$$\frac{\Gamma(a-n)}{\Gamma(a)} = \frac{(-1)^n}{(1-a)_n}, \quad a \neq 0, \pm 1, \pm 2, \dots \quad (3.3)$$

$$(2n)! = 2^{2n} \left(\frac{1}{2}\right)_n n! \quad (3.4)$$

$$(2n+1)! = 2^{2n} \left(\frac{3}{2}\right)_n n! \quad (3.5)(i) \quad \text{Setting } i = j = 0 \text{ in (2.3), we get}$$

$$\frac{1}{\Gamma(a')\Gamma(b)\Gamma(b)} \int_0^\infty \int_0^\infty \int_0^\infty e^{-t_1-t_2-t_3} t_1^{a'-1} t_2^{b-1} t_3^{b-1} \Psi_2^{(3)}(a;d,c,c;xt_1, yt_2, -yt_3) dt_1 dt_2 dt_3$$

$$= X \begin{matrix} 1:2;1 \\ 0:3;1 \end{matrix} \left[ \begin{matrix} a : b, c-b & ; a' \\ - : c, \frac{1}{2}c, \frac{1}{2}c + \frac{1}{2} & ; d \end{matrix} ; \frac{y^2}{4}, x \right]. \quad (3.6)$$

(ii) Setting  $i = 0, j = 1$  in (2.3), we get

$$\frac{1}{\Gamma(a')\Gamma(b)\Gamma(b)} \int_0^\infty \int_0^\infty \int_0^\infty e^{-t_1-t_2-t_3} t_1^{a'-1} t_2^{b-1} t_3^{b-1} \Psi_2^{(3)}(a;d,c,c+1;xt_1, yt_2, -yt_3) dt_1 dt_2 dt_3$$

$$= X \begin{matrix} 1:2;1 \\ 0:3;1 \end{matrix} \left[ \begin{matrix} a : b, c-b+1 & ; a' \\ - : c, \frac{1}{2}c + \frac{1}{2}, \frac{1}{2}c + 1 & ; d \end{matrix} ; \frac{y^2}{4}, x \right]$$

$$+ \frac{aby}{c(c+1)} X \begin{matrix} 1:2;1 \\ 0:3;1 \end{matrix} \left[ \begin{matrix} a+1 : b+1, c-b+1 & ; a' \\ - : c+1, \frac{1}{2}c + 1, \frac{1}{2}c + \frac{3}{2} & ; d \end{matrix} ; \frac{y^2}{4}, x \right] \quad (3.7)$$

(iii) Setting  $i = j = 1$  in (2.3), we get

$$\frac{1}{\Gamma(a')\Gamma(b-1)\Gamma(b)} \int_0^\infty \int_0^\infty \int_0^\infty e^{-t_1-t_2-t_3} t_1^{a'-1} t_2^{b-2} t_3^{b-1} \Psi_2^{(3)}(a;d,c,c+2;xt_1, yt_2, -yt_3) dt_1 dt_2 dt_3$$

$$= X \begin{matrix} 1:2;1 \\ 0:3;1 \end{matrix} \left[ \begin{matrix} a : b, c-b+2 & ; a' \\ - : c+1, \frac{1}{2}c + 1, \frac{1}{2}c + \frac{3}{2} & ; d \end{matrix} ; \frac{y^2}{4}, x \right]$$

$$+ \frac{a(c-2b+2)y}{c(c+2)} X \begin{matrix} 1:2;1 \\ 0:3;1 \end{matrix} \left[ \begin{matrix} a+1 : b, c-b+2 & ; a' \\ - : c+1, \frac{1}{2}c + \frac{3}{2}, \frac{1}{2}c + 2 & ; d \end{matrix} ; \frac{y^2}{4}, x \right]. \quad (3.8)$$

(iv) Setting  $i = 0$  in (2.4), we get

$$\frac{1}{\Gamma(a)} \int_0^\infty e^{-s} s^{a-1} \Phi_2^{(3)}(b,b,c;d;xs,-xs,ys) ds = X \begin{matrix} 1:1;1 \\ 1:0;0 \end{matrix} \left[ \begin{matrix} a:b;c \\ d:-;- \end{matrix} ; x^2, y \right]. \quad (3.9)$$

(v) Setting  $i = -1$  in (2.4), we get

$$\frac{1}{\Gamma(a)} \int_0^\infty e^{-s} s^{a-1} \Phi_2^{(3)}(b+1,b,c;d;xs,-xs,ys) ds$$

$$= X \begin{matrix} 1:1;1 \\ 1:0;0 \end{matrix} \left[ \begin{matrix} a:b+1;c \\ d:-;- \end{matrix} ; x^2, y \right] + \frac{ax}{d} X \begin{matrix} 1:1;1 \\ 1:0;0 \end{matrix} \left[ \begin{matrix} a+1:b+1;c \\ d+1:-;- \end{matrix} ; x^2, y \right]. \quad (3.10)$$

(vi) Setting  $i = 1$  in (2.4), we get

$$\frac{1}{\Gamma(a)} \int_0^\infty e^{-s} s^{a-1} \Phi_2^{(3)}(b-1,b,c;d;xs,-xs,ys) ds$$

$$= X \begin{matrix} 1:1;1 \\ 1:0;0 \end{matrix} \left[ \begin{matrix} a:b;c \\ d:-;- \end{matrix} ; x^2, y \right] - \frac{ax}{d} X \begin{matrix} 1:1;1 \\ 1:0;0 \end{matrix} \left[ \begin{matrix} a+1:b;c \\ d+1:-;- \end{matrix} ; x^2, y \right], \quad (3.11)$$

which for  $x = -x$  and  $b = b+1$  gives the result (3.10).

The other special cases of (2.3) and (2.4) can be obtained by the similar manner .

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## صيغ تكاملية تتضمن الدوال الفوق هندسية ثلاثية المتغيرات $\Psi_2^{(3)}$ و $\Phi_2^{(3)}$

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### الملخص

□

هدف بحثنا هذا هو إثبات صيغ تكاملية عامة متضمنة الدوال الفوق هندسية ثلاثية المتغيرات  $\Psi_2^{(3)}$  و  $\Phi_2^{(3)}$  وذلك بمساعدة صيغتين لدوال لارسيلا  $F_A^{(3)}$  و  $F_D^{(3)}$  ذات الثلاثة المتغيرات والمعطاة في [1] و [2] أيضاً تم عرض بعض التطبيقات لنتائج بحثنا الرئيسية.

**الكلمات المفتاحية:** صيغ تكاملية، الدوال الفوق هندسية، نظرية ديكسون، نظرية كומר.