

On a Generalized βH – Trirecurrent Finsler Space

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Abstract

In this paper, we introduced a Finsler space for which the h – curvature tensor H_{jkh}^i (curvature tensor of Berwald) satisfies the condition

$$\beta_\ell \beta_m \beta_n H_{jkh}^i = c_{\ell mn} H_{jkh}^i + d_{\ell mn} (\delta_k^i g_{jh} - \delta_h^i g_{jk}) - 2y^r b_{mn} \beta_r (\delta_k^i C_{jh\ell} - \delta_h^i C_{jk\ell}) - 2y^r w_{\ell n} \beta_r (\delta_k^i C_{jhm} - \delta_h^i C_{jkm}) - 2y^r \mu_n \beta_\ell \beta_r (\delta_k^i C_{jhm} - \delta_h^i C_{jkm}), H_{jkh}^i = 0,$$

where C_{jkm} is (h) hv – torsion tensor, $\beta_\ell \beta_m \beta_n$ is Berwald's covariant differential operator of the third order with respect to x^n , x^m and x^ℓ , successively, $\beta_\ell \beta_r$ is Berwald's covariant differential operator of the second order with respect to x^ℓ and x^r , successively, β_r is Berwald's covariant differential operator of the first order with respect to x^r , $c_{\ell mn}$ and $d_{\ell mn}$ are non – zero covariant tensors field of third order, b_{mn} and $w_{\ell n}$ are non – zero covariant tensors field of second order and μ_ℓ is non – zero covariant vector field. We called this space a *generalized βH – trirecurrent space*. The aim of this paper is to develop some properties of a generalized βH – trirecurrent space by obtaining Berwald's covariant derivative of the third order for the (h)v – torsion tensor H_{kh}^i and the deviation tensor H_k^i , the curvature vector H_k and the scalar curvature H are investigated.

Key words: Finsler space, generalized βH – trirecurrent space, Ricci tensor.

1. Introduction

Pandey P.N., Saxena S.S. and Goswami A. [3] introduced and studied a generalized H– recurrent Finsler space. F.Y.A.Qasem [4] introduced and discussed generalized H– birecurrent curvature tensor and W.H.A. Hadi [1] studied the generalized – birecurrent for some tensors and studied some special spaces in this space.

Let F_n be an n – dimensional Finsler spaces equipped with the metric function F satisfies conditions [5], the tensor C_{ijk} is positively homogeneous of degree -1 in y^i and symmetric in all its indices and is called (h)hv – torsion tensor [2]. According to Euler's theorem on homogeneous functions, this tensor satisfies the following:

$$(1.1) \quad C_{ijk} y^i = C_{jki} y^i = C_{kij} y^i = 0.$$

Berwald's covariant derivative $\mathcal{B}_k T_j^i$ of an arbitrary tensor field T_j^i with respect to x^k is given by [5]

$$(1.2) \quad \mathcal{B}_k T_j^i := \partial_k T_j^i - (\partial_r T_j^i) G_k^r + T_j^r G_{rk}^i - T_r^i G_{jk}^r.$$

Berwald's covariant derivative of y^i vanish as identically [5], i.e.

$$(1.3) \quad \mathcal{B}_k y^i = 0.$$

In view of (1.2), the second covariant derivative of an arbitrary vector field X^i with respect to x^h in the sense of Berwald [5],

$$(1.4) \quad \mathcal{B}_h \mathcal{B}_k X^i = \partial_k (\mathcal{B}_h X^i) - (\partial_s \mathcal{B}_h X^i) G_k^s - (\mathcal{B}_r X^i) G_{hk}^r + (\mathcal{B}_h X^r) G_{rk}^i.$$

Using (1.4) and taking skew – symmetric part, with respect to the indices k and h, we get the commutation formula for Berwald's covariant differentiation as follows [5]:

$$(1.5) \quad \mathcal{B}_h \mathcal{B}_k X^i - \mathcal{B}_k \mathcal{B}_h X^i = X^r H_{rkh}^i - (\dot{\partial}_r X^i) H_{kh}^r$$

where

$$(1.6) \quad (a) \quad H_{jkh}^i := \partial_h G_{jk}^i + G_{jk}^r G_{rh}^i + G_{rjh}^i G_k^r - h/k \quad \text{and} \quad (b) \quad H_{kh}^i := \partial_h G_k^i + G_k^r G_{rh}^i - h/k.$$

Remark 1.1. $-h/k$ means the subtraction from the former term by interchanging the indices h and k.

The tensors H_{jkh}^i and H_{kh}^i , as defined above, are called *h – curvature tensor* (*h – curvature tensor of Berwald*) and *h(hv) – torsion tensor* are positively homogeneous of degree zero and one in y^i , respectively.

Berwald constructed the tensors H_{jkh}^i and H_{kh}^i from the tensor H_k^i called by him as *deviation tensor*, according to

$$(1.7) \quad H_h^i := 2\partial_h G^i + \partial_s G_h^i y^s + 2G_{hs}^i G^s - G_s^i G_h^s.$$

The h(hv) – torsion tensor and the deviation tensor satisfy the following [5]:

$$(1.8) \quad (a) \quad H_{jkh}^i y^j = H_{kh}^i, \quad (b) \quad H_{jk}^i y^j = H_k^i \quad \text{and} \quad (c) \quad \dot{\partial}_j H_k^i = H_{jk}^i.$$

The H – Ricci tensor, the curvature vector and the scalar curvature satisfy the following [5]:

$$(1.9) \quad (a) \quad H_{jki}^i = H_{jk}, \quad (b) \quad H_{ki}^i = H_k \quad \text{and} \quad (c) \quad H_i^i = H.$$

2. Generalized βH – Trirecurrent Space

A Finsler space for which Berwald curvature tensor H_{jkh}^i satisfies the generalized recurrence property, i.e. characterized by the equation [3]

$$(2.1) \quad \beta_n H_{jkh}^i = \lambda_n H_{jkh}^i + \mu_n (\delta_k^i g_{jh} - \delta_h^i g_{jk}), \quad H_{jkh}^i = 0,$$

where β_n is the differential operator with respect to x^n in the sense of Berwald, λ_n and μ_n are non – zero covariant vectors field and isvcalled *the recurrence vectors field*, such space known as *generalized H – recurrent Finsler space*.

Taking the β – covariant derivative for (2.1) with respect to x^m , we get

$$\beta_m \beta_n H_{jkh}^i = (\beta_m \lambda_n) H_{jkh}^i + \lambda_n (\beta_m H_{jkh}^i) + (\beta_m \mu_n) (\delta_k^i g_{jh} - \delta_h^i g_{jk}) + \mu_n \beta_m (\delta_k^i g_{jh} - \delta_h^i g_{jk}).$$

In view of the equation (2.1), we can write the above equation as

$$\beta_m \beta_n H_{jkh}^i = (\beta_m \lambda_n + \lambda_n \lambda_m) H_{jkh}^i + (\lambda_n \mu_m + \beta_m \mu_n) (\delta_k^i g_{jh} - \delta_h^i g_{jk}) - 2y^r \mu_n \beta_r (\delta_k^i C_{jhm} - \delta_h^i C_{jkm})$$

which can be written as [4]

$$(2.2) \quad \beta_m \beta_n H_{jkh}^i = a_{mn} H_{jkh}^i + b_{mn} (\delta_k^i g_{jh} - \delta_h^i g_{jk}) - 2y^r \mu_n \beta_r (\delta_k^i C_{jhm} - \delta_h^i C_{jkm}),$$

where $a_{mn} = \beta_m \lambda_n + \lambda_n \lambda_m$ and $b_{mn} = \lambda_n \mu_m + \beta_m \mu_n$ are non – zero covariant tensors field of second order and $\beta_m \beta_n$ is the differential operator with respect to x^n and x^m , successively, such space known as *generalized βH – birecurrent space*.

Remark 2.1. The expression β – covariant derivative stands the covariant derivative in the sense of Berwald.

Taking the β – covariant derivative for (2.2) with respect to x^ℓ , we get

$$\begin{aligned} \beta_\ell \beta_m \beta_n H_{jkh}^i &= (\beta_\ell a_{mn}) H_{jkh}^i + a_{mn} (\beta_\ell H_{jkh}^i) + (\beta_\ell b_{mn}) (\delta_k^i g_{jh} - \delta_h^i g_{jk}) \\ &+ b_{mn} \beta_\ell (\delta_k^i g_{jh} - \delta_h^i g_{jk}) - 2y^r (\beta_\ell \mu_n) \beta_r (\delta_k^i C_{jhm} - \delta_h^i C_{jkm}) \\ &- 2y^r \mu_n \beta_\ell \beta_r (\delta_k^i C_{jhm} - \delta_h^i C_{jkm}). \end{aligned}$$

In view of the equation (2.1), we can write the above equation as

$$\begin{aligned} \beta_\ell \beta_m \beta_n H_{jkh}^i &= (\beta_\ell a_{mn} + a_{mn} \lambda_\ell) H_{jkh}^i + (\beta_\ell b_{mn} + a_{mn} \mu_\ell) (\delta_k^i g_{jh} - \delta_h^i g_{jk}) \\ &\quad - 2y^r b_{mn} \beta_r (\delta_k^i C_{jhl} - \delta_h^i C_{jkl}) - 2y^r (\beta_\ell \mu_n) \beta_r (\delta_k^i g_{jh} - \delta_h^i g_{jk}) \\ &\quad - 2y^r \mu_n \beta_\ell \beta_r (\delta_k^i C_{jhm} - \delta_h^i C_{jkm}) \end{aligned}$$

which can be written as

$$(2.3) \quad \beta_\ell \beta_m \beta_n H_{jkh}^i = c_{\ell mn} H_{jkh}^i + d_{\ell mn} (\delta_k^i g_{jh} - \delta_h^i g_{jk}) - 2y^r b_{mn} \beta_r (\delta_k^i C_{jhl} - \delta_h^i C_{jkl}) - 2y^r w_{\ell n} \beta_r (\delta_k^i C_{jhm} - \delta_h^i C_{jkm}),$$

where $c_{\ell mn} = \beta_\ell a_{mn} + a_{mn} \lambda_\ell$ and $d_{\ell mn} = \beta_\ell b_{mn} + a_{mn} \mu_\ell$ are non – zero covariant tensors field of third order $w_{\ell n} = \beta_\ell \mu_n$ is non – zero covariant tensor field of second order and $\beta_\ell \beta_m \beta_n$ is the differential operator with respect to x^n, x^m and x^ℓ , successively.

Definition 2.1. A Finsler space F_n for which Berwald curvature tensor H_{jkh}^i satisfies the condition (2.3) and called a *generalized $\beta H -$ trirecurrent space* and the tensor a *generalized $\beta -$ trirecurrent*. We shall denote them briefly as $G \beta H - TR$ and $G\beta - TR$, respectively.

Now, transvecting the condition (2.3) by y^j , in view of (1.3), and by using (1.8a) and (1.1), we get

$$(2.4) \quad \beta_\ell \beta_m \beta_n H_{kh}^i = c_{\ell mn} H_{kh}^i + d_{\ell mn} (\delta_k^i y_h - \delta_h^i y_k)$$

Transvecting (2.4) by y^h , in view of (1.3) and by using (1.8b), we get

$$(2.5) \quad \beta_\ell \beta_m \beta_n H_k^i = c_{\ell mn} H_k^i + d_{\ell mn} (\delta_k^i F^2 - y^i y_k).$$

Thus, we may conclude

Theorem 2.1. *In $G \beta H - TR F_n$, Berwald's covariant derivative of the third order for the $h(v) -$ torsion tensor H_{kh}^i and the deviation tensor H_k^i given by the equations (2.4) and (2.5), respectively.*

Contracting the indices i and k in the condition (2.3) and using (1.9a), we get

$$(2.6) \quad \beta_\ell \beta_m \beta_n H_{jh} = c_{\ell mn} H_{jh} + d_{\ell mn} (n - 1) g_{jh} - 2b_{mn} (n - 1) \beta_r C_{jhl} - 2y^r w_{\ell n} (n - 1) \beta_r C_{jhm} - 2y^r \mu_n (n - 1) \beta_\ell \beta_r C_{jhm}.$$

Thus, we may conclude

Theorem 2.2. *In $G \beta H - TR F_n$, Berwald's covariant derivative of the third order for the $H -$ Ricci tensor H_{jk} given by the equation (2.6).*

Contracting the indices i and k in the equations (2.4) and (2.5), using (1.9b) and (1.9c), we get

$$(2.7) \quad \beta_\ell \beta_m \beta_n H_h = c_{\ell mn} H_h + d_{\ell mn} (n - 1) y_h$$

and

$$(2.8) \quad \beta_\ell \beta_m \beta_n H = c_{\ell mn} H + d_{\ell mn} F^2.$$

The equations (2.7) and (2.8) show that the curvature vector H_h and the scalar curvature H can't vanish because the vanishing of them would imply $d_{\ell mn} = 0$, a contradiction.

Thus, we may conclude

Theorem 2.3. *In $G \beta H - TR F_n$, the curvature vector H_h and the curvature scalar H are non – vanishing.*

We know that [5]

$$(2.9) \quad H_{ikh}^i = H_{hk} - H_{kh}.$$

Taking the $\beta -$ covariant derivative of the third order with respect to x^n, x^m and x^ℓ , successively, for the equation (2.9), we get

$$\beta_\ell \beta_m \beta_n H_{ikh}^i = \beta_\ell \beta_m \beta_n H_{hk} - \beta_\ell \beta_m \beta_n H_{kh} .$$

By using (2.6), the above equation can be written

$$\begin{aligned} \beta_\ell \beta_m \beta_n H_{ikh}^i &= c_{\ell mn} H_{hk} + d_{\ell mn} (n - 1) g_{hk} - 2y^r b_{mn} (n - 1) \beta_r C_{hkl} - \\ 2y^r w_{\ell n} (n - 1) \beta_r C_{hkm} & - 2y^r \mu_n (n - 1) \beta_\ell \beta_r C_{hkm} - [c_{\ell mn} H_{kh} + d_{\ell mn} (n - 1) g_{kh} - \\ 2y^r b_{mn} (n - 1) \beta_r C_{khl} & - 2y^r w_{\ell n} (n - 1) \beta_r C_{khm} - 2y^r \mu_n (n - 1) \beta_\ell \beta_r C_{khm}]. \end{aligned}$$

By using the symmetric property of the h(hv) – torsion tensor C_{jkh} in it's all indies , the above equation can be written as

$$\beta_\ell \beta_m \beta_n H_{ikh}^i = c_{\ell mn} (H_{hk} - H_{kh}).$$

By using (2.9) in the above equation , we get

$$\beta_\ell \beta_m \beta_n (H_{hk} - H_{kh}) = c_{\ell mn} (H_{hk} - H_{kh}).$$

Thus , we may conclude

Theorem 2.4. *In $G \beta H - TR F_n$, the tensor $(H_{hk} - H_{kh})$ behaves as trirecurrent .*

We know that [5]

$$(2.10) \quad H_{kh}^i = \frac{1}{3} (\partial_k H_h^i - \partial_h H_k^i).$$

Taking the $\beta -$ covariant derivative of third order with respect to x^n, x^m and x^ℓ , successively, for the equation (2.10), we get

$$\beta_\ell \beta_m \beta_n H_{kh}^i = \frac{1}{3} \beta_\ell \beta_m \beta_n (\partial_k H_h^i - \partial_h H_k^i) .$$

In view of (1.8c) , the above equation can be written as

$$\beta_\ell \beta_m \beta_n H_{kh}^i = \frac{1}{3} (\beta_\ell \beta_m \beta_n H_{kh}^i - \beta_\ell \beta_m \beta_n H_{hk}^i) .$$

By using (2.4) in the above equation, we get

$$\beta_\ell \beta_m \beta_n H_{kh}^i = \frac{1}{3} [c_{\ell mn} H_{kh}^i + d_{\ell mn} (\delta_k^i y_h - \delta_h^i y_k) - c_{\ell mn} H_{hk}^i - d_{\ell mn} (\delta_h^i y_k - \delta_k^i y_h)]$$

which can be written as

$$\beta_\ell \beta_m \beta_n H_{kh}^i = \frac{1}{3} c_{\ell mn} (H_{kh}^i - H_{hk}^i) + \frac{2}{3} d_{\ell mn} (\delta_k^i y_h - \delta_h^i y_k)$$

or

$$\beta_\ell \beta_m \beta_n H_{kh}^i = v_{\ell mn} (H_{kh}^i - H_{hk}^i) + w_{\ell mn} (\delta_k^i y_h - \delta_h^i y_k),$$

where $v_{\ell mn} = \frac{1}{3} c_{\ell mn}$ and $w_{\ell mn} = \frac{2}{3} d_{\ell mn}$.

Thus , we may conclude

Theorem 2.5. *In $G \beta H - TR F_n$, Berwald torsion tensor H_{kh}^i is generalized – trirecurrent tensor.*

References

1. **Hadi, W.H.A.** (2016) : *Study of certain types of generalized birecurrent in Finsler space*, Ph.D. Thesis, University of Aden, (Aden) (Yemen).
2. **Matsumoto, M.** (1971) : On h – Isotropic and C^h – recurrent Finsler , J. Math. Kyoto Univ., 11 ,1–9 .
3. **Pandey, P.N., Saxena, S. and Goswani, A.** (2011) : On a generalized H – recurrent space, Journal of International Academy of Physical Science, Vol. 15, 201–211.
4. **Qasem, F.Y.A.** (2016) : On generalized H – birecurrent Finsler space, International Journal of Mathematics and its Applications, Vol. 4, Issue 2, 51 – 57.
5. **Rund, H.** (1981) : *The differential geometry of Finsler space*, Springer –Verlag, Berlin Gottingen – Heidelberg, (1959); 2nd edit. (in Russian), Nauka, (Moscow) .

حول تعميم فضاء فنسلر βH – ثلاثي المعاودة

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المُلخَص

تم تقديم تعريف هذا الفضاء من خلال اشتقاق المعادلة المميزة له لتعميم فضاء βH – ثنائي المعاودة وأطلقنا عليه الفضاء المعمم βH – ثلاثي المعاودة ورمزنا له بالرمز $G\beta H - TRF_n$ ، كما تم التوصل إلى بعض المبرهنات المختلفة بهذا الخصوص. وقد تم إثبات أن الموتر H – ريشي H_{jn} والمتجه التقوسي H_n هي موترات لا تنتهي في هذا الفضاء.

الكلمات المفتاحية: فضاء فنسلر، مشتقة متحدة الاختلاف بمفهوم بروالد من المرتبة الثالثة – موتر H – ريشي .